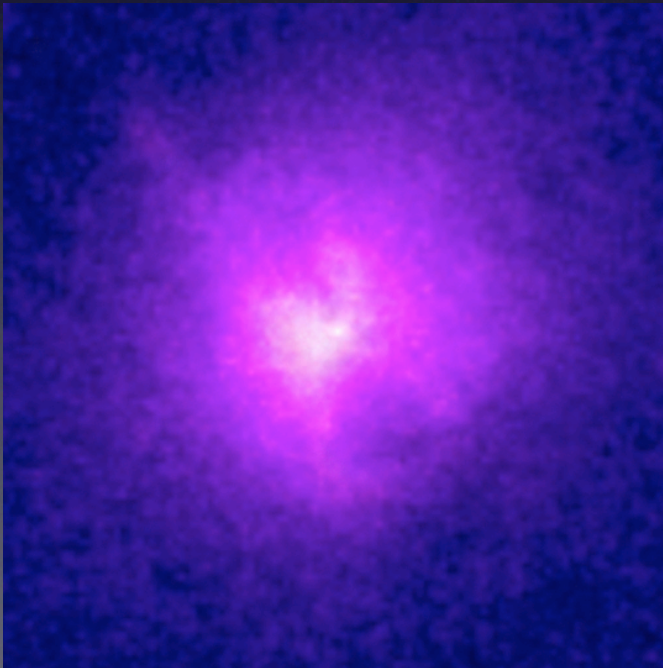


# **New Forms of Convection in Galaxy Cluster Plasmas (i.e., how do galaxy clusters boil?)**

Eliot Quataert (UC Berkeley)

Hydra A w/ Chandra



in collaboration with

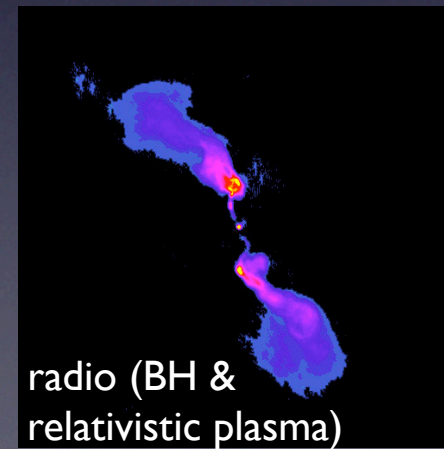
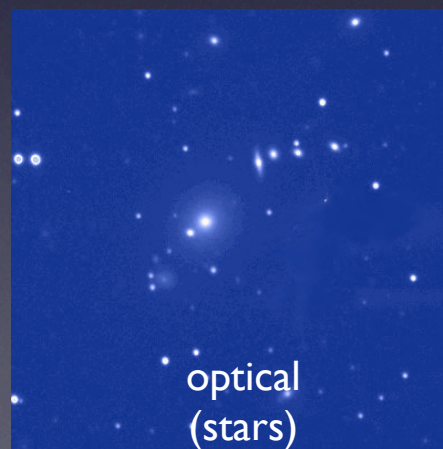
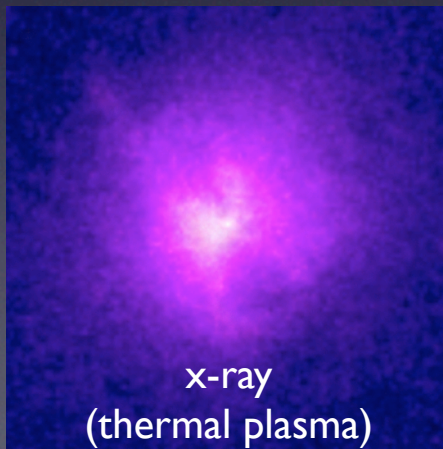
Ian Parrish  
Prateek Sharma

# Overview

- Hot Plasma in Clusters of Galaxies
- Hydrodynamic Convection ('normal' convection; e.g., the sun)
- Convection induced by Anisotropic Thermal Conduction
  - new convective instabilities: the "MTI" & "HBI"
- Implications for Clusters
  - incl. interaction btw. thermal plasma & cosmic rays from an AGN

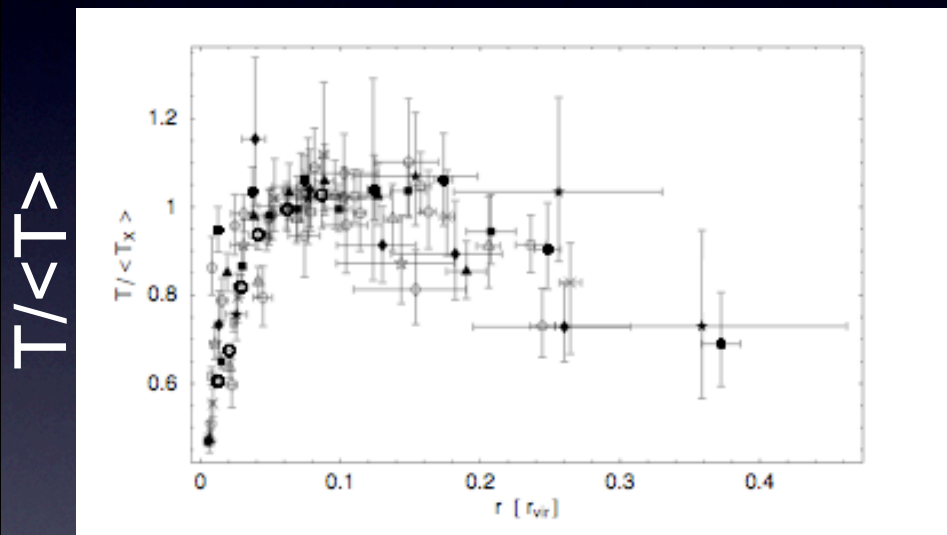
# Clusters of Galaxies

- largest gravitationally bound objects:  $M_{\text{vir}} \sim 10^{14-15} M_{\odot}$   
 $R_{\text{vir}} \sim 1-3 \text{ Mpc}$ 
  - $\sim 84\%$  dark matter;  $\sim 14\%$  **plasma**;  $\sim 2\%$  stars
  - on exponential tail of the mass function: useful cosmological probe
  - host the most massive galaxies ( $\sim 10^{12} M_{\odot}$ ) and black holes ( $\sim 10^{9-10} M_{\odot}$ )



# Hot Plasma in Clusters

cluster temperature profiles



Piffaretti et al. 2005

Radius ( $R_{vir}$ )

$$L_x \sim 10^{43-46} \text{ erg s}^{-1}$$

$$n \sim 10^{-4}-1 \text{ cm}^{-3}$$

$$T \sim 1-15 \text{ keV}$$

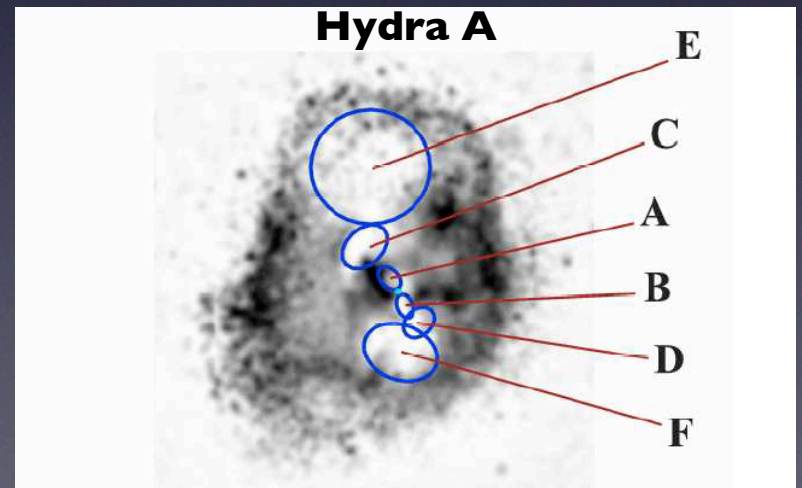
large electron mean free path:

$$\ell_e \simeq 2 \left( \frac{T}{3 \text{ keV}} \right)^2 \left( \frac{n}{0.01 \text{ cm}^{-3}} \right)^{-1} \text{ kpc}$$

→ **thermal conduction**  
**important**

# “Cool Core” Clusters

- in at least  $\sim 50\%$  of clusters,  $t_{\text{cool}} < \text{Hubble time}$  for  $r \lesssim 100 \text{ kpc}$
- absent a heat source:  $\dot{M}_{\text{cool}} \sim 100\text{-}1000 M_{\odot} \text{ yr}^{-1}$ 
  - not observed:  $\dot{M}_{\text{star}} \lesssim 0.01 \dot{M}_{\text{cool}}; T_{\text{min}} \sim 1/3 T_{\text{vir}}$
- $\rightarrow$  a heat source balances radiative cooling
  - $\sim$  spherically out to  $\sim 100\text{s kpc}$
- proposed sources of heating include
  - a central (radio loud) AGN  $\longleftrightarrow$
  - thermal conduction from large R
  - .....



# Hydrodynamic Convection

- Schwarzschild criterion for convection:  **$ds/dz < 0$**
- Motions slow & adiabatic: **pressure equil,  $s \sim \text{const}$**

low entropy ( $s$ )

↓  
gravity  
high  $s$

$$\begin{matrix} \rho_f \\ s_{bg} \\ p'_{bg} \end{matrix}$$

$$\begin{matrix} \rho_i = \rho_{bg} \\ s_i = s_{bg} \\ p_i = p_{bg} \end{matrix}$$

background fluid

$$s'_{bg} \quad \rho'_{bg} \quad p'_{bg}$$

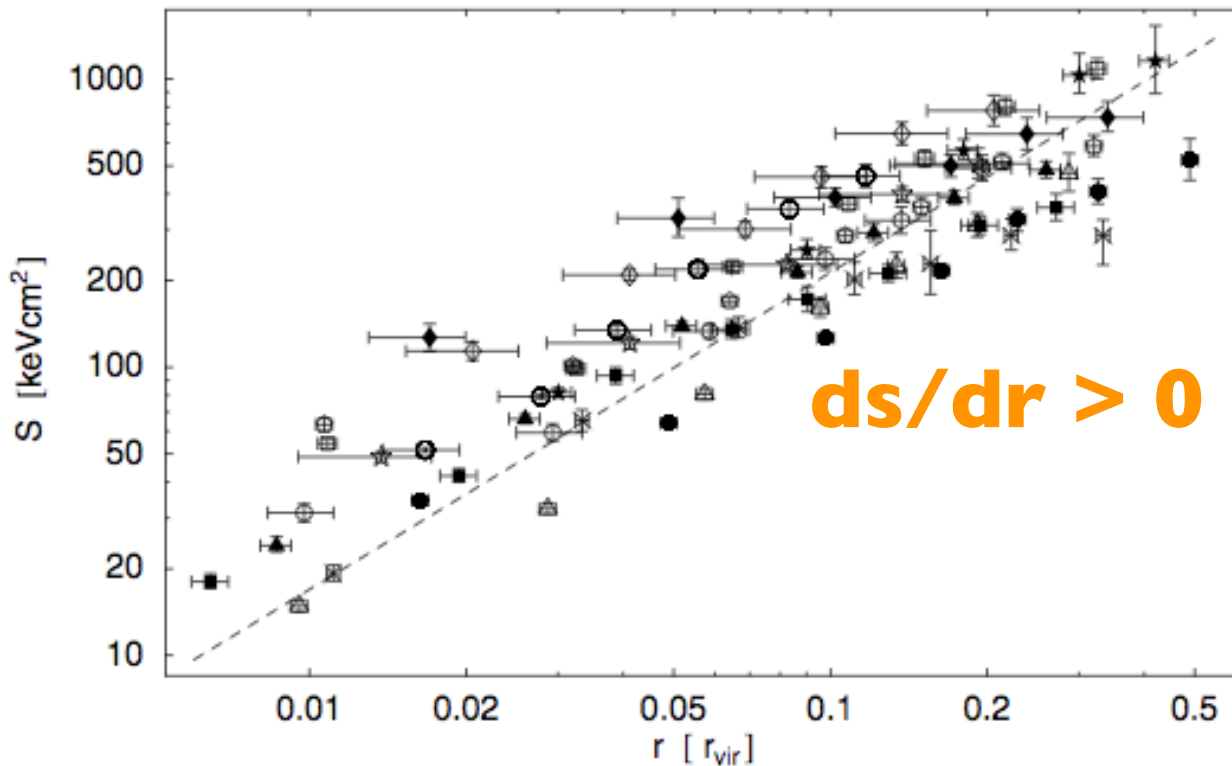
$$s(p, \rho) \propto \ln[p/\rho^\gamma]$$

$$\text{if } ds/dz < 0 \rightarrow \rho_f < \rho'_{bg}$$

**convectively unstable**

# Cluster Entropy Profiles

Entropy



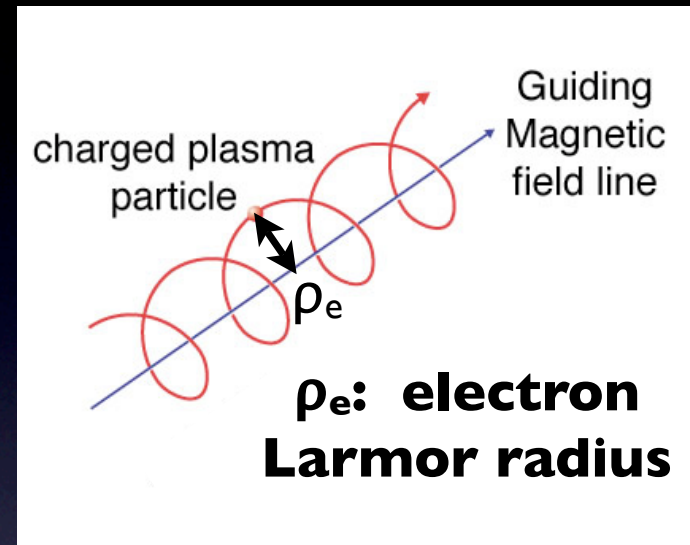
Radius ( $R_{\text{vir}}$ )

Schwarzschild criterion  $\rightarrow$  clusters are **stable**

## **Anisotropic** Thermal Conduction in Cluster Plasmas

electron mean free path:

$$l_e \simeq 2 \left( \frac{T}{3 \text{ keV}} \right)^2 \left( \frac{n}{0.01 \text{ cm}^{-3}} \right)^{-1} \text{ kpc}$$



$$\frac{l_e}{\rho_e} \sim 10^{14} \left( \frac{B}{10^{-6} \text{ G}} \right) \left( \frac{n}{0.01 \text{ cm}^{-3}} \right)^{-1} \left( \frac{T}{3 \text{ keV}} \right)^{-3/2}$$

$l_e \gg \rho_e \Rightarrow$  heat transport is **anisotropic** (primarily along B)



# The Magnetothermal Instability (MTI)

Balbus 2000, 2001; Parrish & Stone 2005, 2007; Quataert 2008; Sharma, Quataert, & Stone 2008

cold



g  
hot

$$\begin{aligned} \rho_i &= \rho_{bg} \\ T_i &= T_{bg} \\ p_i &= p_{bg} \end{aligned}$$

thermal conduction time  
 $\ll$  buoyancy time

$$\begin{aligned} p_f &= p'_{bg} \\ T_f &= T_i \end{aligned}$$

$$\begin{aligned} T_f &> T'_{bg} \\ \rho_f &< \rho'_{bg} \end{aligned}$$

**convectively  
unstable  
( $dT/dz < 0$ )**

weak B-field  
no dynamical effect;  
only channels heat flow

growth time  
 $\sim$  dyn. time

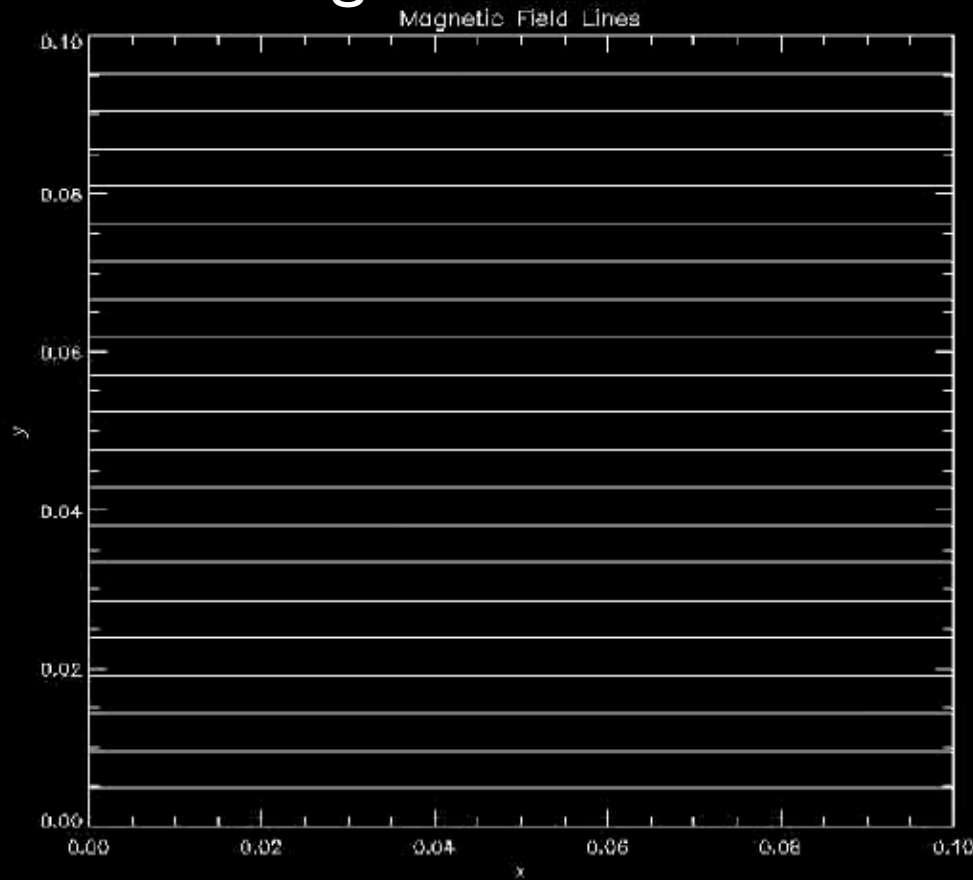
# The MTI

magnetic field lines

cold



hot



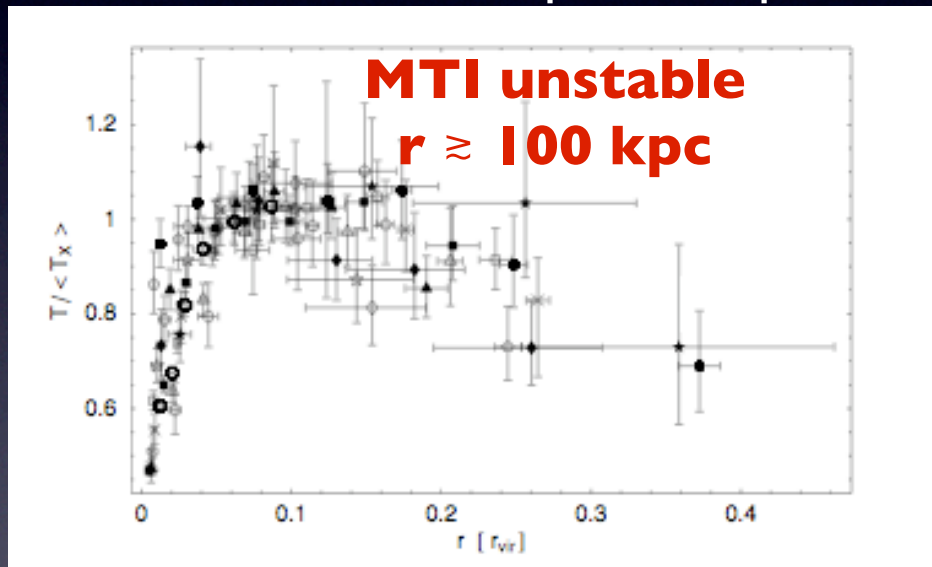
2D simulation from Ian Parrish

instability saturates by rearranging  
magnetic field lines

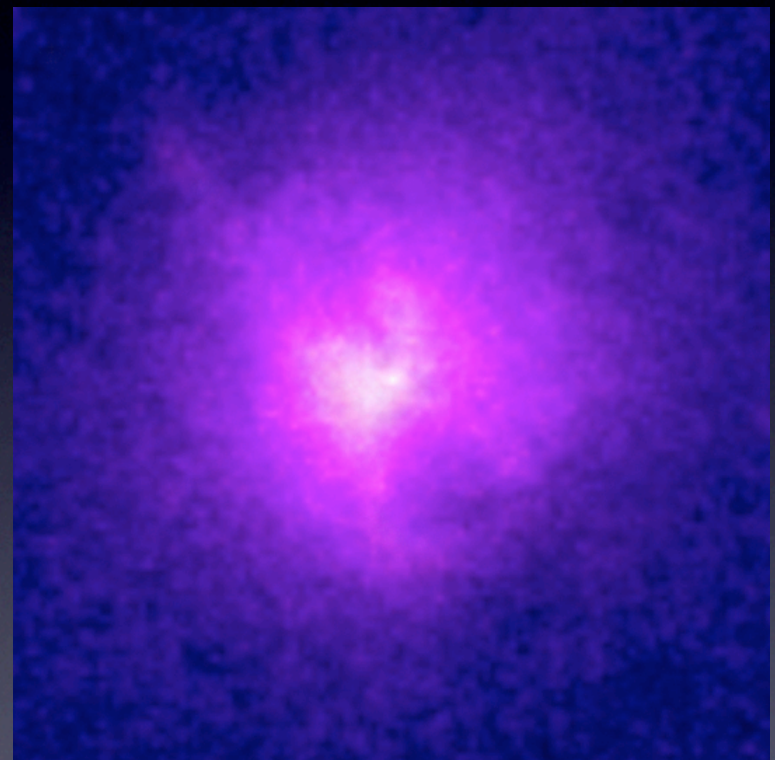
# The MTI in Clusters

cool core cluster temperature profile

$T_{\text{core}}/T_{\text{edge}}$



Radius ( $R_{\text{vir}}$ )



Piffaretti et al. 2005

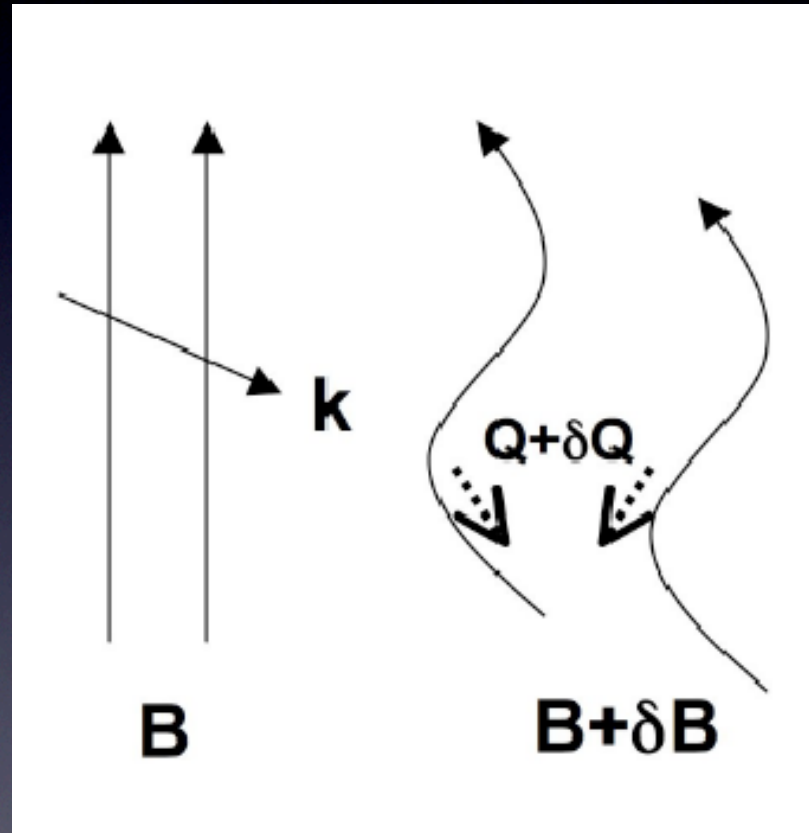
$\sim 200$  kpc

# The Heat Flux-Driven Buoyancy Instability (HBI)

Quataert 2008; Parrish & Quataert 2008

hot  
↓  
cold  
 $g, Q_z$   
**heat flux**

↑  
weak  
 $B$



converging &  
diverging  
heat flux

⇒

conductive  
heating &  
cooling

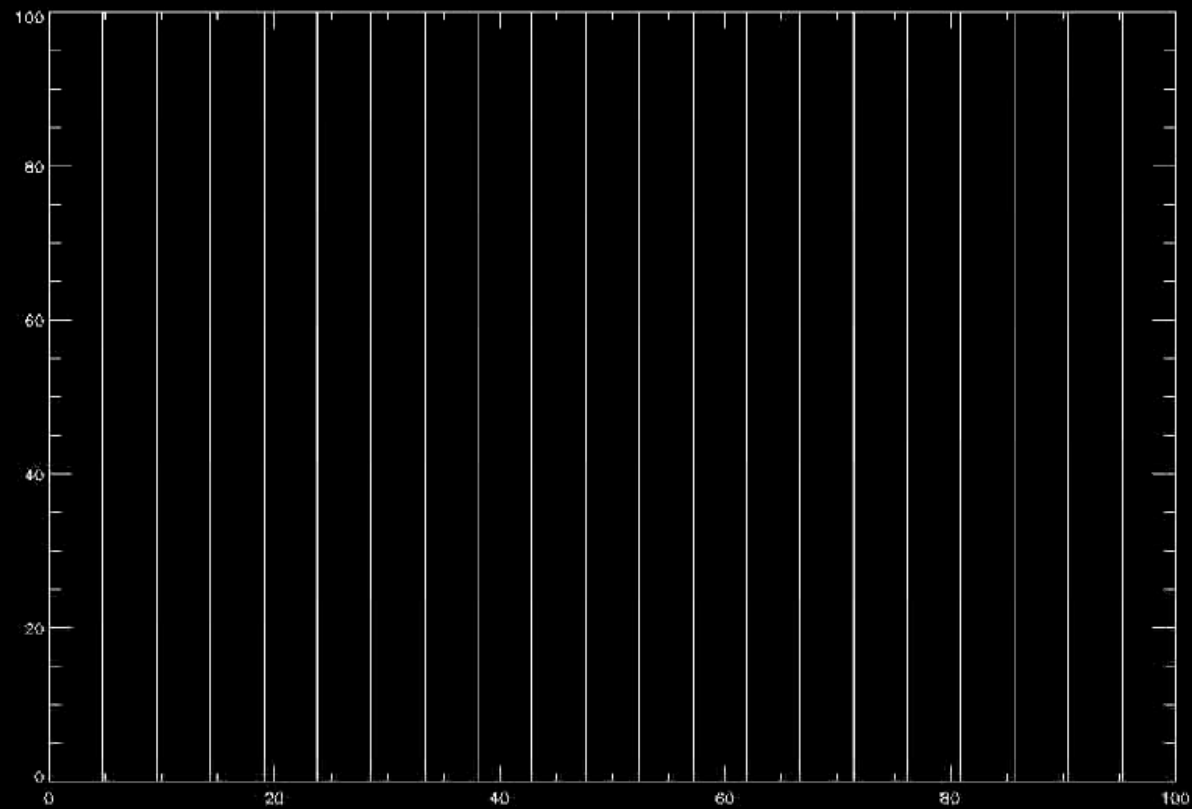
for  $dT/dz > 0$   
upwardly displaced  
fluid is heated  
& rises

**convectively  
unstable**

# The HBI

hot  
↓  
cold  
g,  $Q_z$   
bg heat flux

magnetic field lines

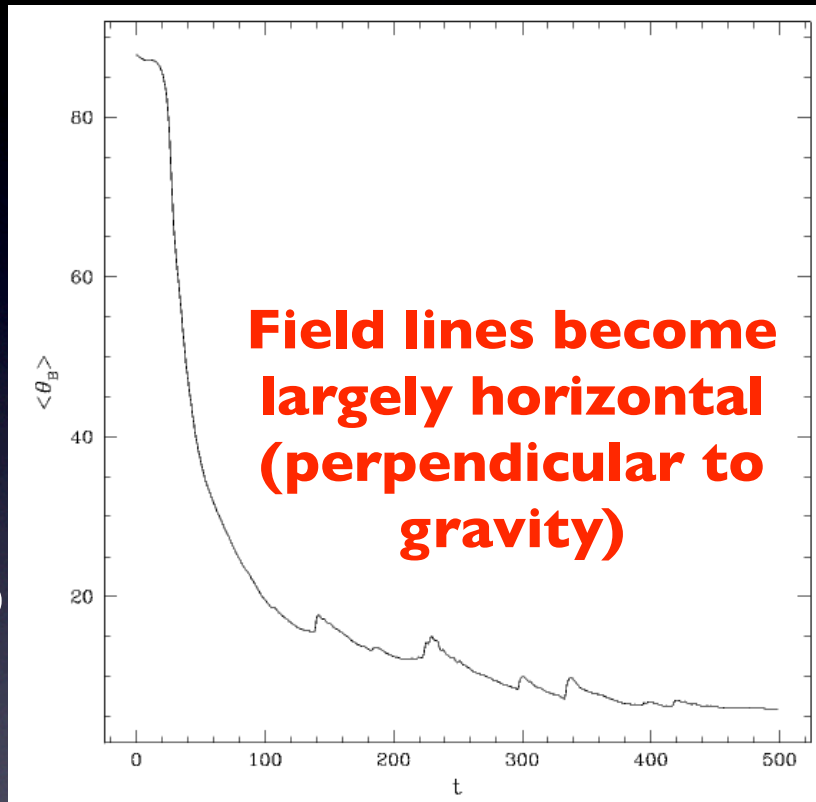


Parrish & Quataert 2008

HBI saturates by rearranging  
the magnetic field lines

# Nonlinear Evolution: HBI

Angle wrt Horizontal

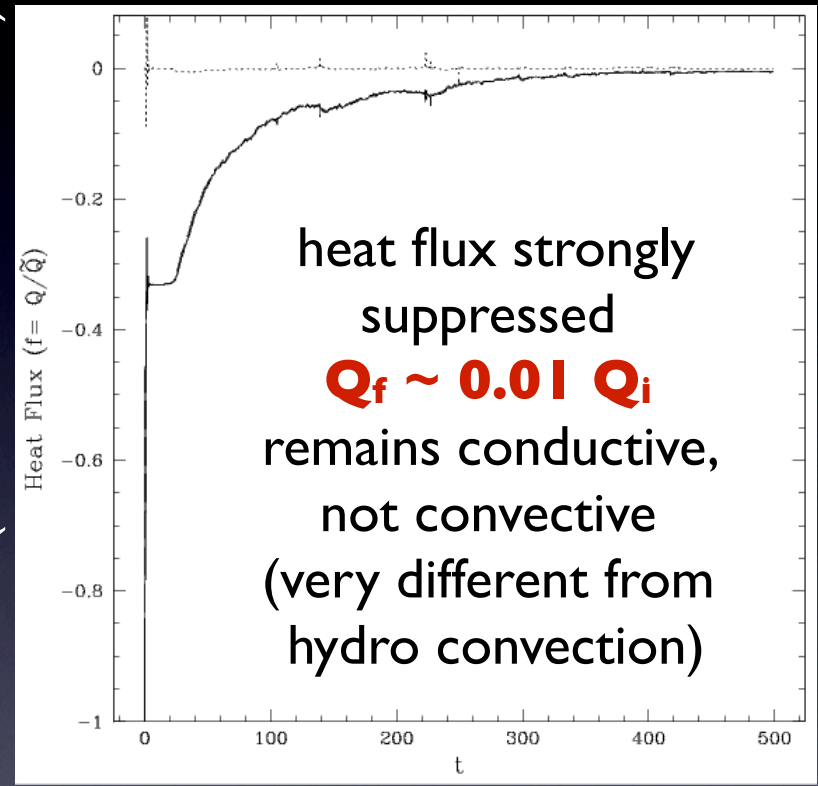


Partish & Quataert 2008

time (dynamical time)

B-field energy  
amplified by  $\sim 100$

Heat Flux (relative to initial value)

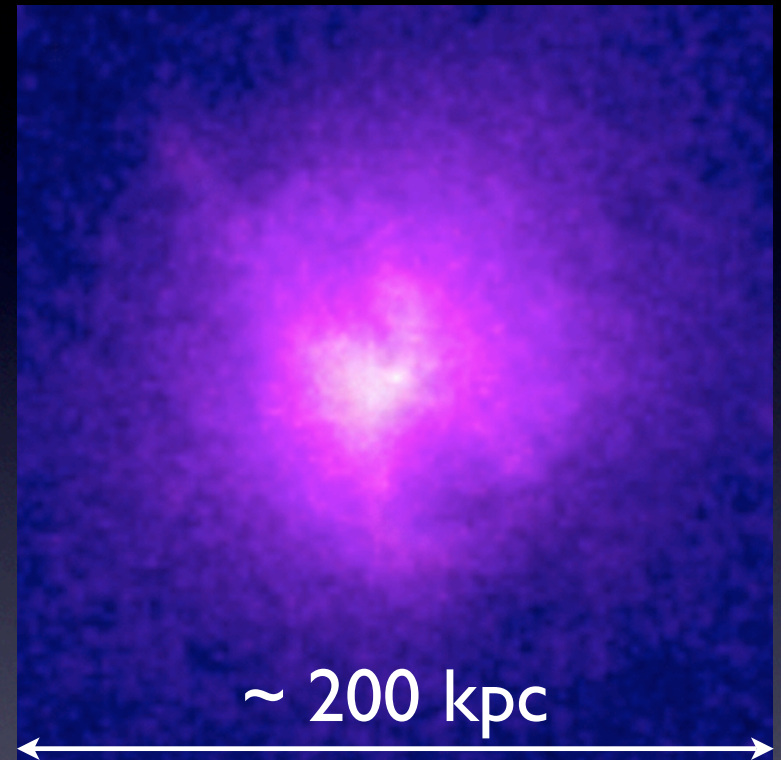
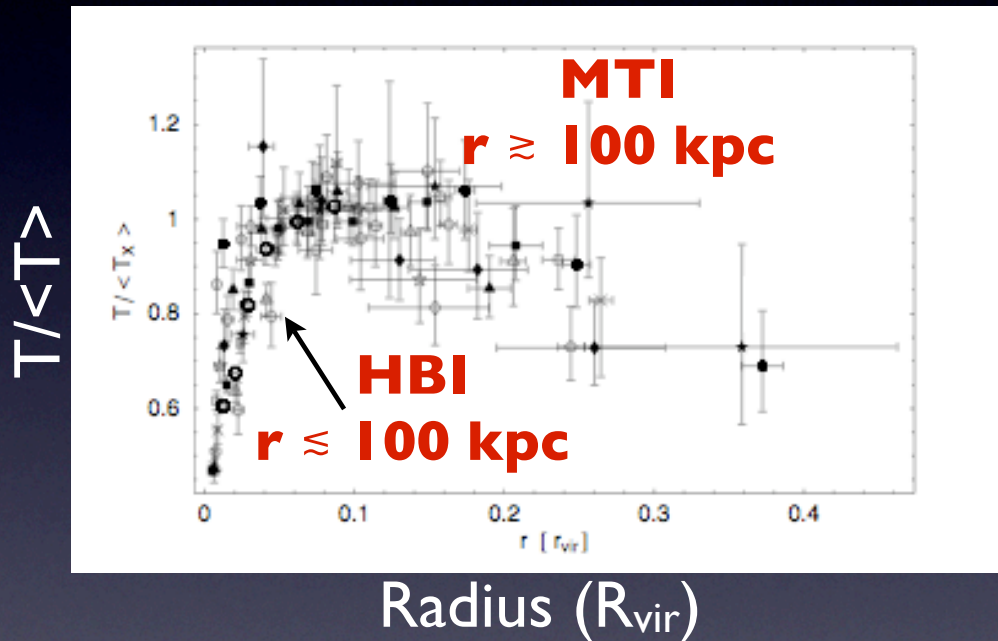


time (dynamical time)

**Local 3D Simulations**  
**initially weak B; no cooling**

# The MTI & HBI in Clusters

cool core cluster temperature profile



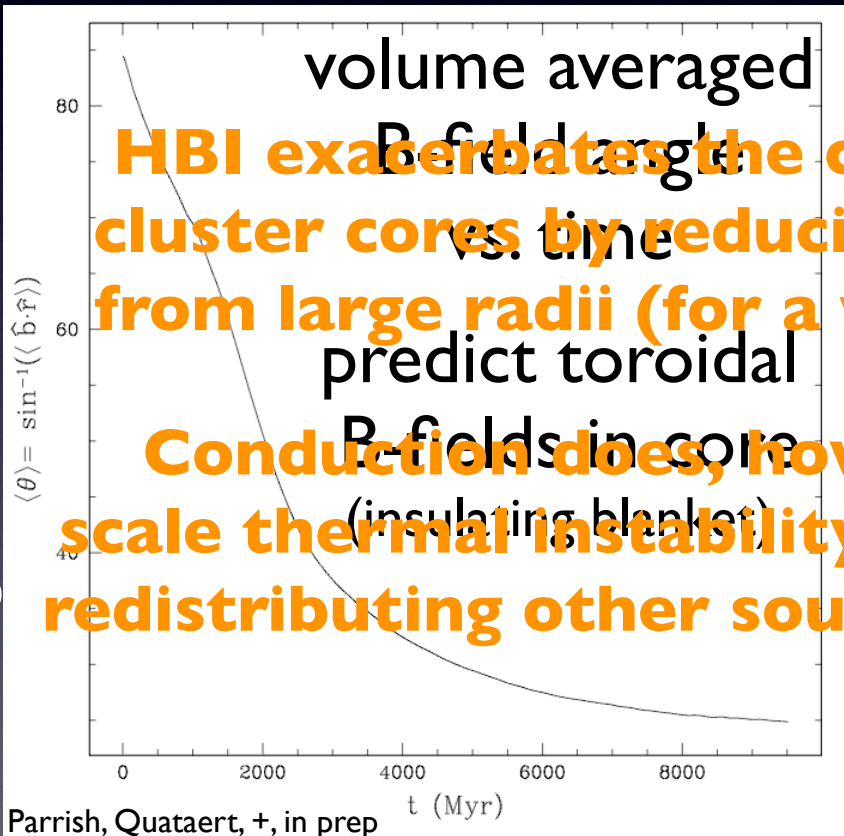
**The Entire Cluster is Convectively Unstable!**

1. strong B (e.g., solar corona) or
2. isotropic heat transport  $\gg$  anisotropic heat transport (e.g., solar interior)

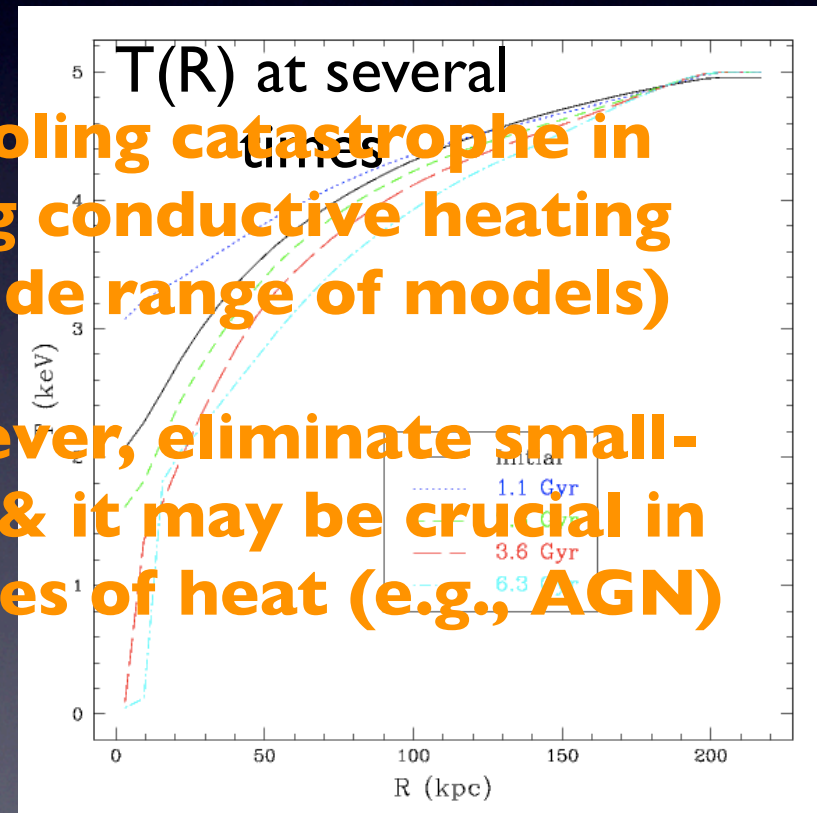
# Global Cluster Simulations

- 3D **w/ cooling & anisotropic conduction** (Athena)
  - non-cosmological: isolated cluster core ( $\approx 200$  kpc)
  - conductive flux is not a “free parameter”; depends on dynamics!

Angle wrt Horizontal



time (Myr)



Radius (kpc)

**HBI exacerbates the cooling catastrophe in cluster cores by reducing conductive heating from large radii (for a wide range of models)**

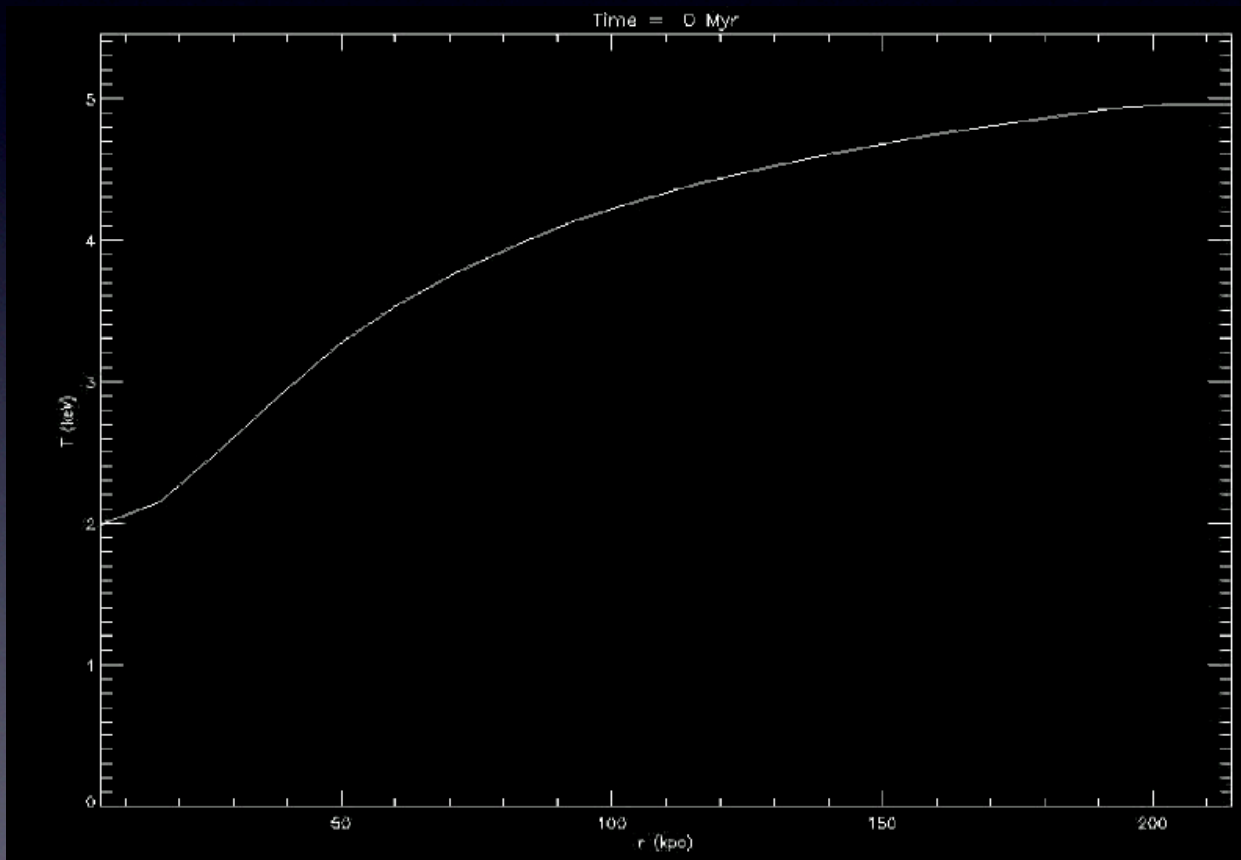
**Conduction does, however, eliminate small-scale thermal instability; & it may be crucial in redistributing other sources of heat (e.g., AGN)**



# Global Cluster Simulations

- 3D **w/ cooling & anisotropic conduction** (Athena)
- *artificial* source of heating to balance cooling at  $< 20$  kpc (“AGN”)

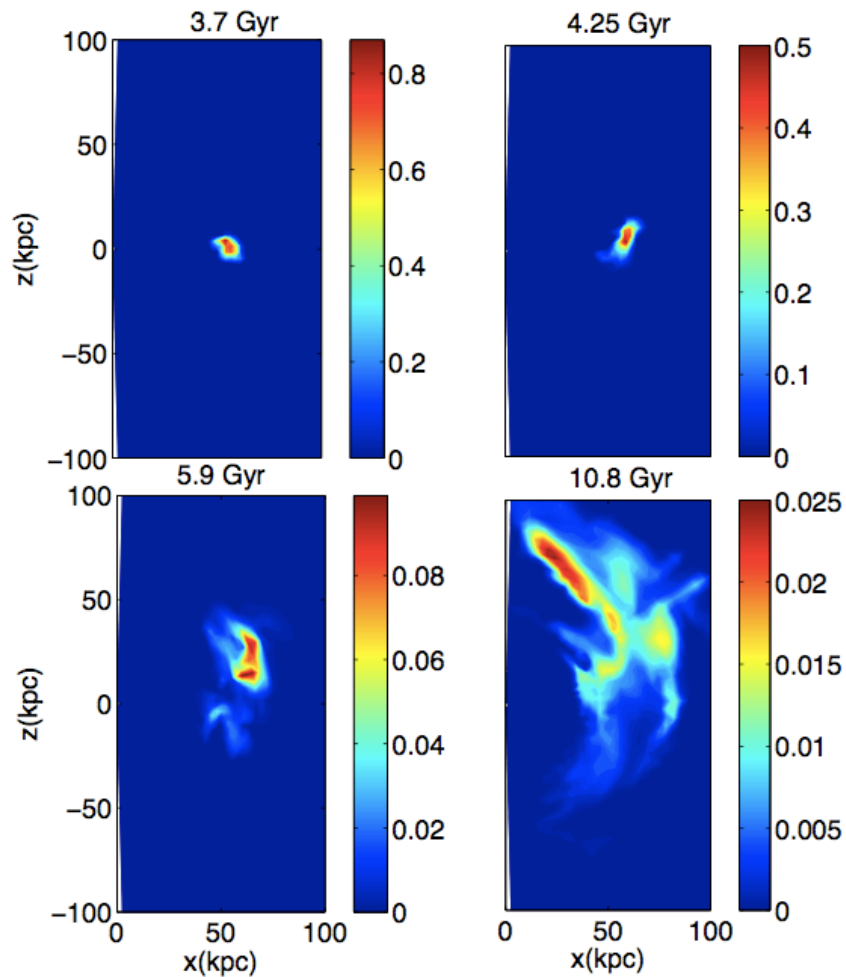
Temperature (keV)



Radius (kpc)

**Stable for  
~ 10 Gyr**

# HBI-induced Turbulence

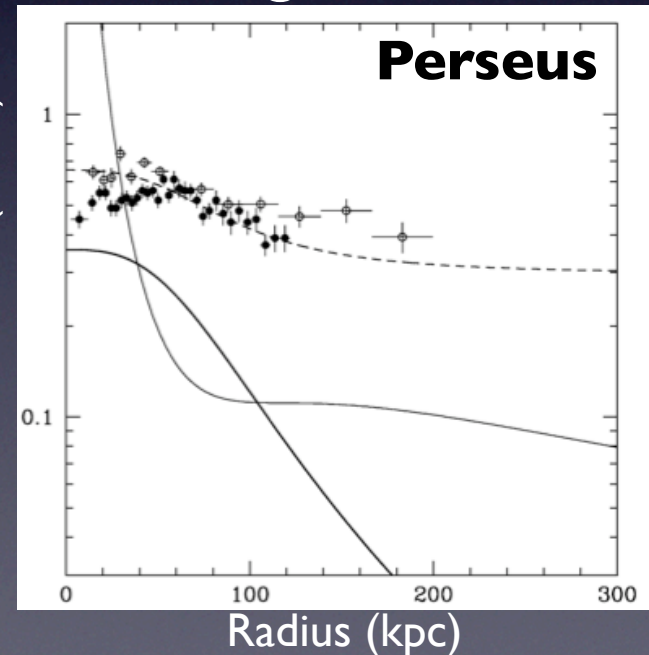


Density of Passive Scalar  
Linear Color Scale  
(red/blue = high/low density)

$v_{\text{turb}} \sim 0.01-0.1 c_s$   
detectable w/ next  
generation x-ray  
calorimeters

mixing of metals

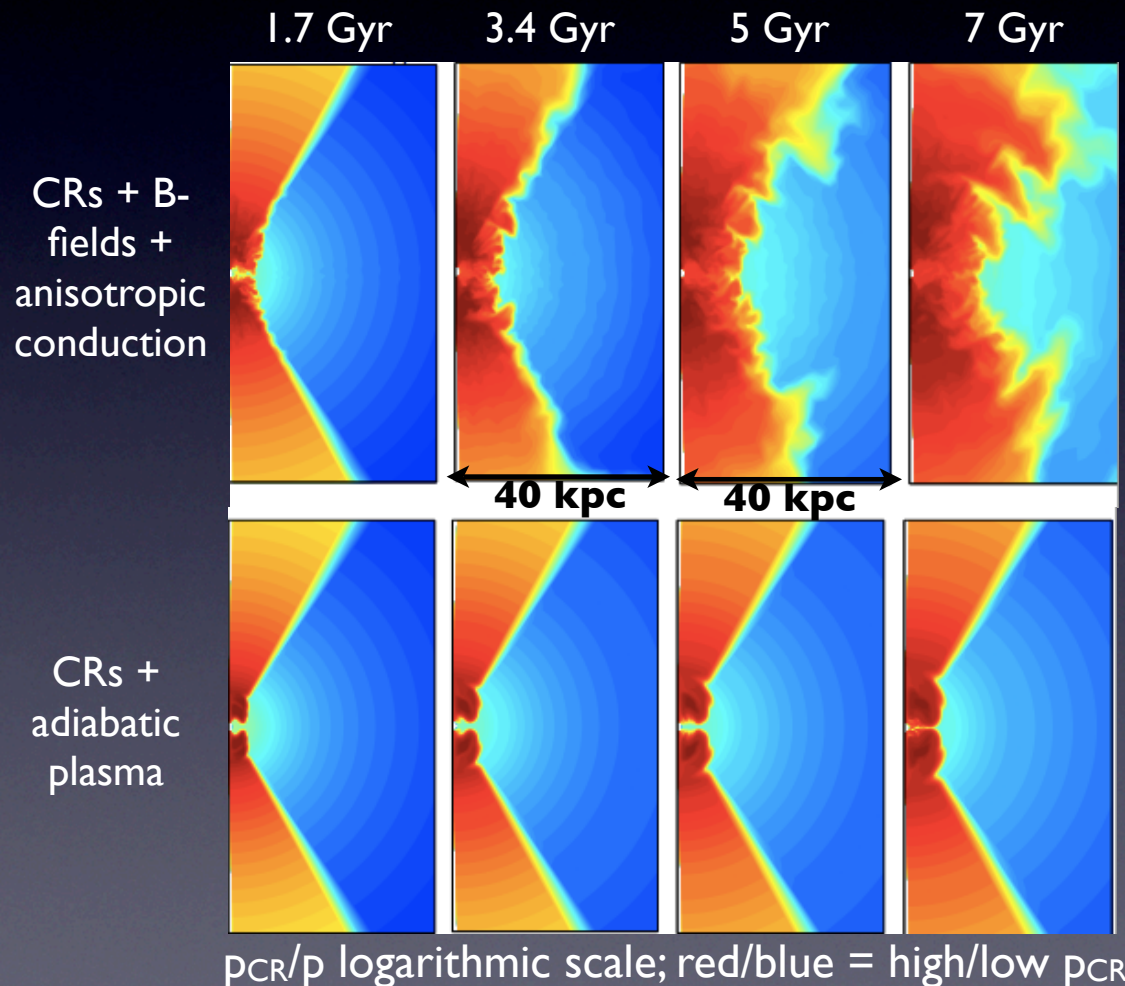
Fe abundance (solar)



Robusco et al. 2005

# Effects on CR Mixing

AGN heating is the most promising mechanism balancing cooling; but precise physical mechanism & how it couples **throughout** the cluster core unclear



“real” cluster plasma:  
buoyantly unstable &  
easier to mix CRs

adiabatic plasma:  
buoyantly stable &  
harder to mix CRs

# Summary

- Understanding the thermal history of galaxy cluster cores is a key to understanding the process of massive galaxy/BH formation
- Recent Surprises: the plasma throughout a galaxy cluster is convectively unstable (MTI & HBI)!
  - key role of anisotropic thermal conduction (accept no substitutes)
  - HBI inhibits conductive heating of cluster cores
- The Future: interplay between AGN heating, cosmic rays, and realistic cluster thermodynamics